

THE EFFECT OF ND-NANOCUSTER INTERACTION ON DE-EXCITATION OF ND-DOPED SILICON-RICH SILICON OXIDE

Se-Young Seo and Jung H. Shin

Silicon Photonics Laboratory
Department of Physics, Korea Advanced Institute of
Science and Technology (KAIST), 373-1 ~Kusung-
dong, Yuseong-gu, Taejeon, Korea

Outline

- Rare-earth doped nc-Si
- Excitation/de-excitation of RE in nc-Si
- Factors suppress thermal quenching (energy mismatch and others?)
- Probing thermal quenching with Nd-doping
- Nd excitation via exciton recombination of nc-Si
- Weak thermal quenching process in SRSO:Nd
- The lack of backtransfer due to weak couplings
- Implications

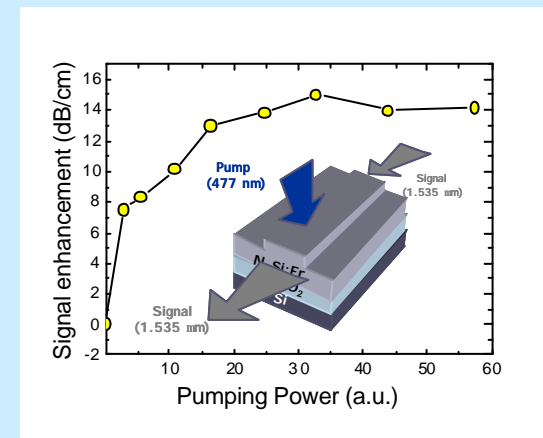
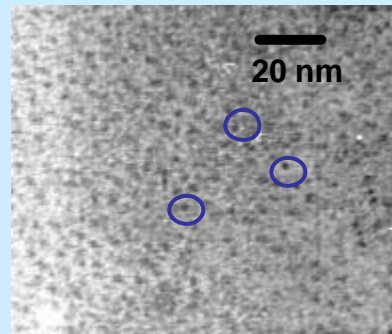
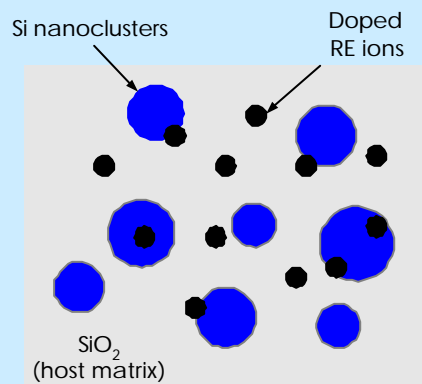
Rare-earth doped Si nanoclusters

▪ Rare-earth doped Si

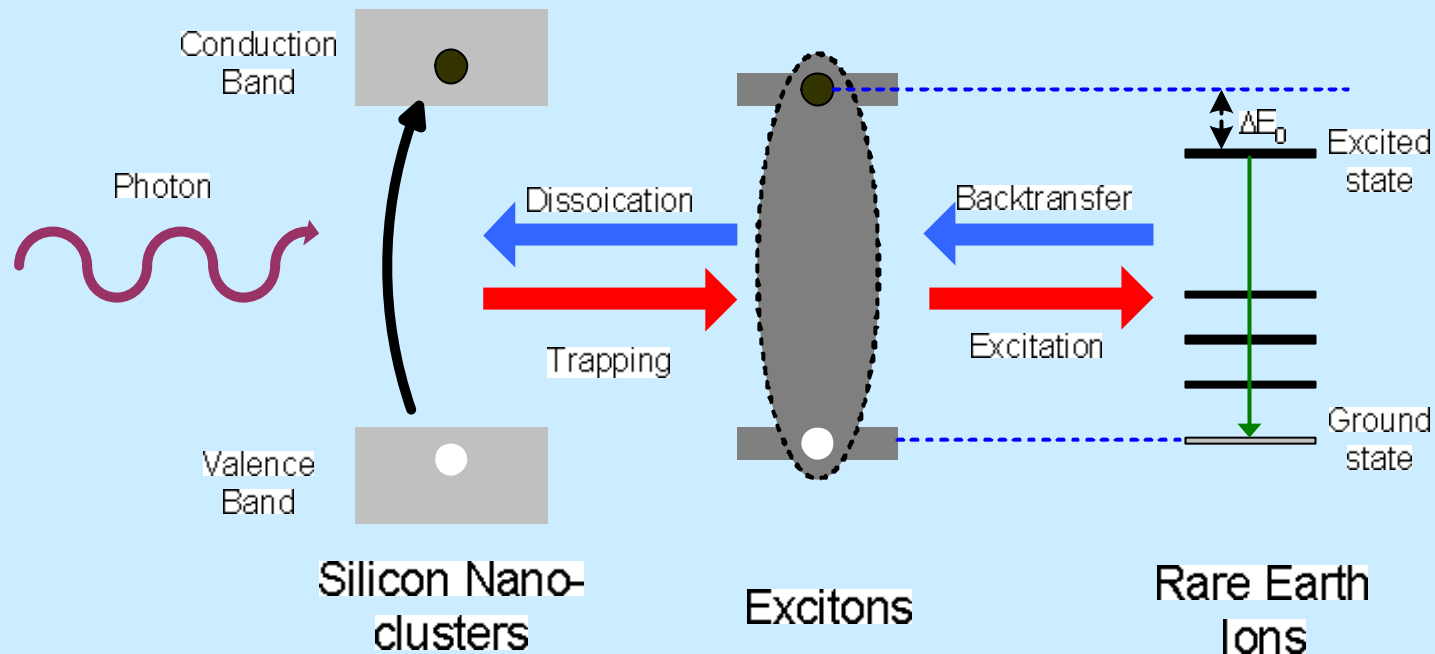
- Si : well-developed infrastructure: from knowledge to technology
- Rare-earth: stable and intrinsic luminescence by 4f transition independent of host and temperature
- ➔ Rare-earth doped Si: realization of silicon compatible micro-photonics

▪ Silicon rich silicon oxide (SRSO):RE

- SRSO: Si nanoclusters embedded in SiO_2
- Efficient RE luminescence environment
(Wide bandgap, high excitation and luminescence efficiency)
- ➔ 1.54 μm optical gain from nc-Si:Er (Han *et al.* APL 2001)



Excitation/De-excitation of RE in nc-Si



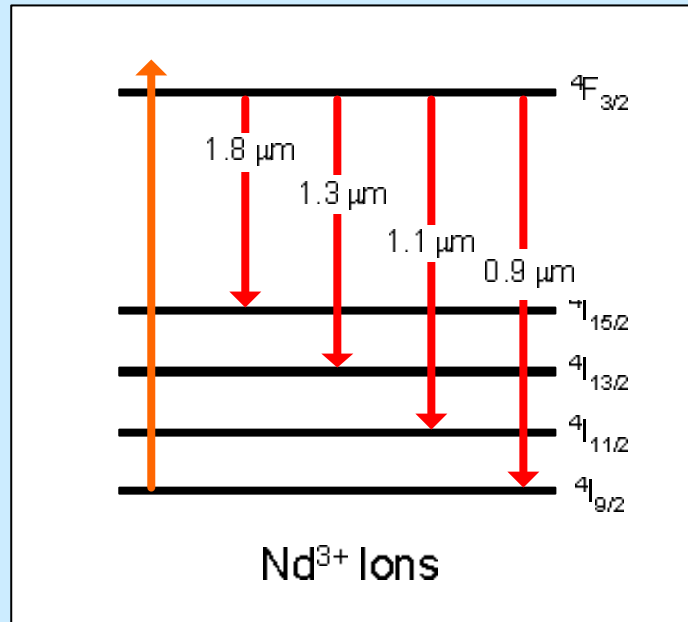
▪ Exciton mediated excitation

- Similar to the RE excitation of bulk semiconductors (Schmitt-Rink PRL 1991)
- Backtransfer: non-radiative decay of RE with help of phonons
 - Dominant thermal quenching mechanism
 - Plagues many RE-doped bulk semiconductors

Suppression of thermal quenching

- **Wide bandgap:** increase of $E_0 \rightarrow$ increase of energy barriers (e.g. Favennec *et al.* MRS Proc. 1993)
 - Wide bandgap host: SiC, GaN
 - Quantum confinement effect of nano-clustering of narrow bandgap
 - \rightarrow Bandgap widening with decrease of nc size
 - \rightarrow The lack of thermal quenching for nc-Si (e.g. Seo *et al.* APL 1999)
- **Micro-structure: bandgap is not the whole picture**
 - Slight separation of RE away from nc-Si can also suppress thermal quenching
 - \rightarrow Different quenching behavior between Er-immersed and implanted porous Si (Wang *et al.* J. Luminescence)
 - \rightarrow Suppression of thermal quenching by buffer SiO₂ layer between nc-Si and SiO₂:Er (Shin *et al.* APL)
- **What factors other than bandgap contribute to suppression of backtransfer?**

Probing of backtransfer with Nd-doping



▪ Nd³⁺ luminescence

- 0.9 μm ($4F_{3/2} \rightarrow 4I_{9/2}$) and 1.1 μm ($\rightarrow 4I_{11/2}$)
 - within c-Si and nc-Si luminescence range
 - Sensitive probe of the interaction between nc-Si and RE ions
- 0.9 μm luminescence can be detectable with Si photodiode
 - Integrating light emission and detection capabilities into single Si chip

Experiment

▪ **Film Fabrication**

- ECR-PECVD with co-sputtering of Nd³⁺ target
- Stoichiometric study : Si (34-50 at. %)
- 1 mm thick with 0.14 at. % Nd

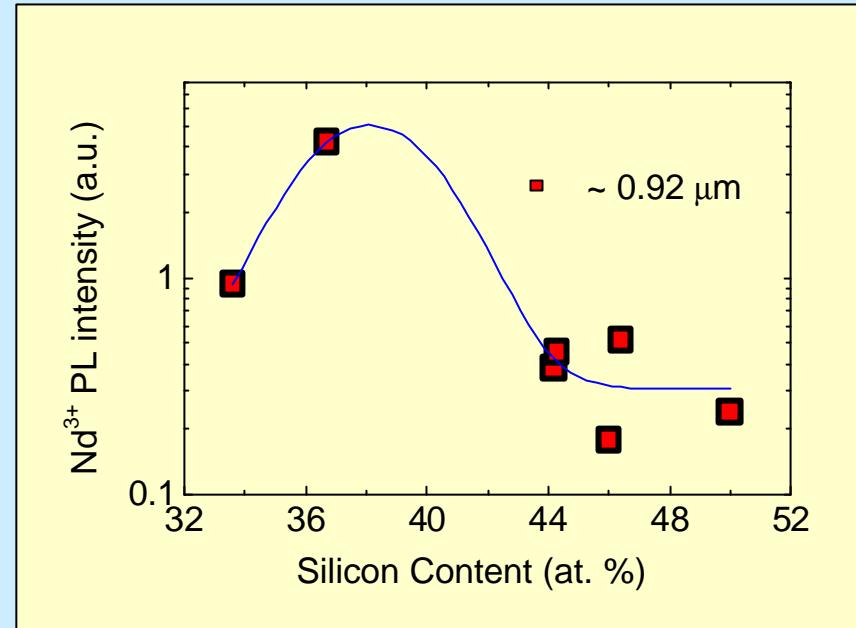
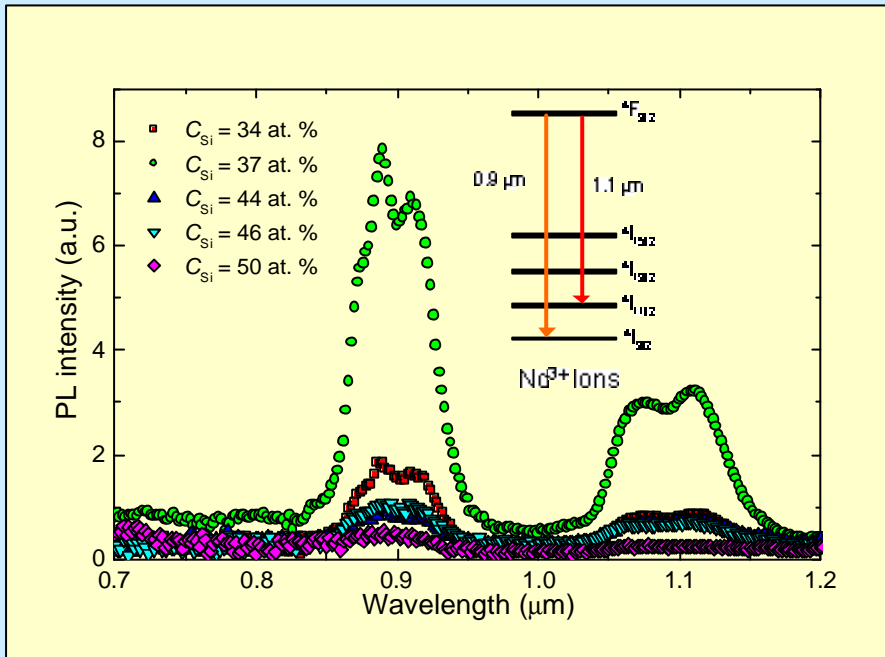
▪ **Post Treatment**

- Rapid thermal anneal: 950 C for 5 min. for clustering of nc-Si

▪ **Luminescence measurement**

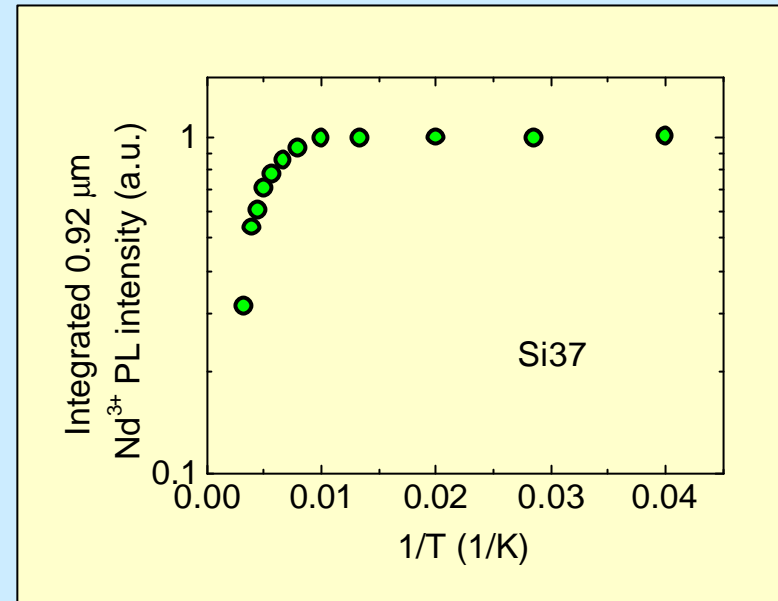
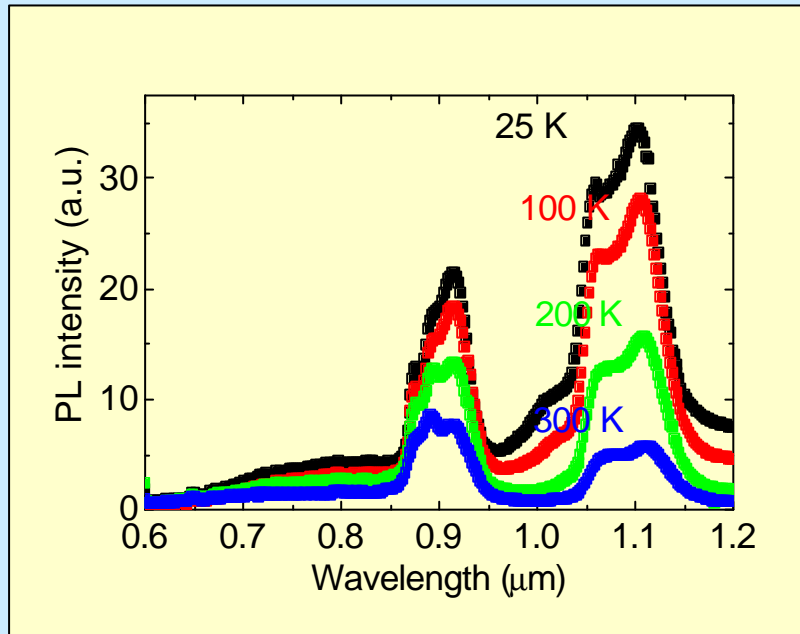
- Detectors: Si photodiode + InGaAs photodiode
- Pumped using **488 nm** line of Ar ion laser: not resonant with Nd³⁺ levels
- Recording of luminescence decay traces with InGaAs(Cs) PMT and digitizing oscilloscope

The size of nc-Si : crucial for Nd³⁺ excitation



- **Clear Nd³⁺ luminescence at room-temperature**
 - Both 0.92 μm and 1.1 μm luminescence observed
- **Typical RE PL intensity dependence on Si content**
 - Nc-Si with > 44 at. % Si : too large not to have sufficient bandgap for Nd³⁺ excitation
 - The crucial role of quantum confinement effect of nc-Si for Nd-excitation
→ Nd³⁺ excitation via exciton recombination

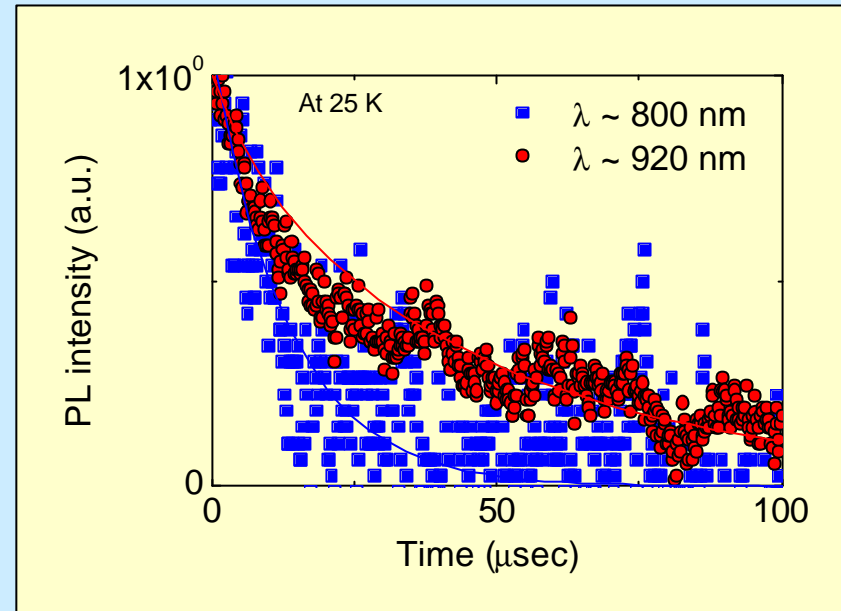
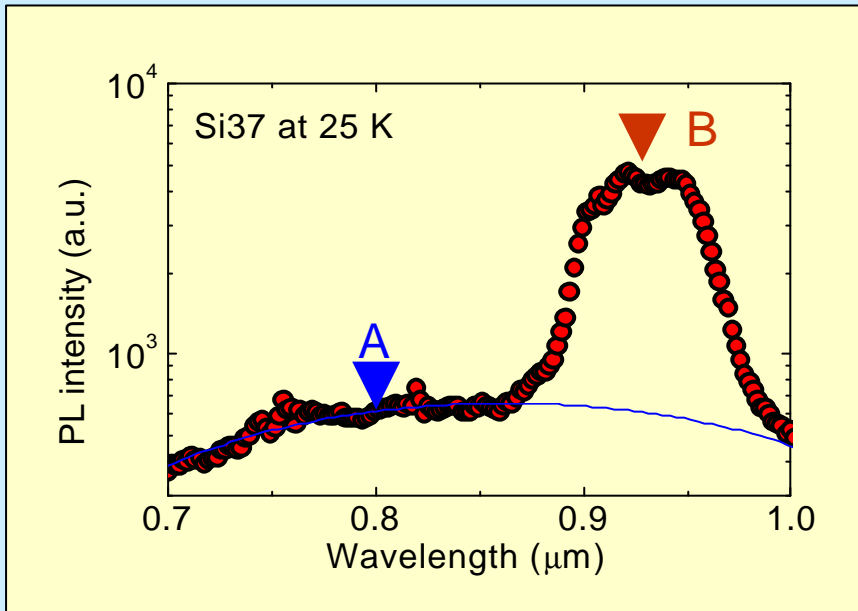
Temperature dependence of Nd³⁺ PL



- Intrinsic nc-Si luminescence at $\sim 0.8 \mu\text{m}$
- For $0.92 \mu\text{m}$ luminescence
 - The superimposition of Nd³⁺ PL on the tail of intrinsic nc-Si luminescence
 - 3 folds decrease of Nd³⁺ PL intensity for Si37

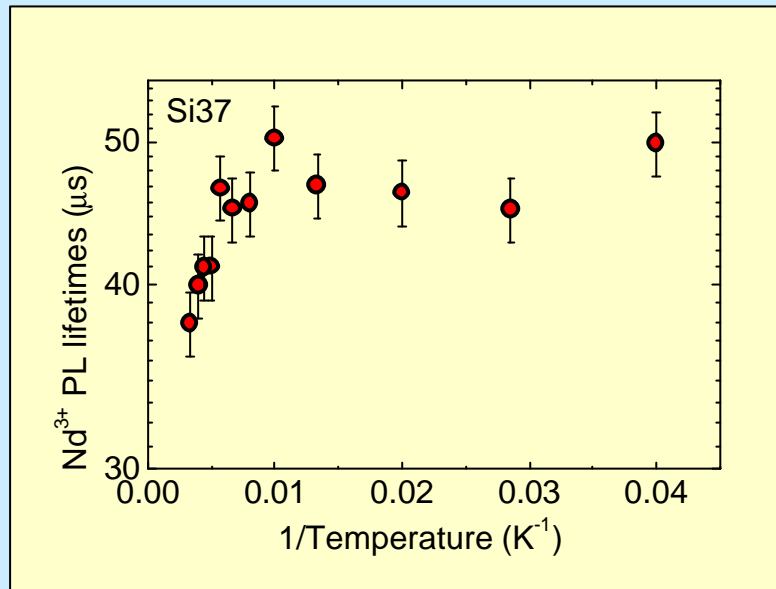
The determination of Nd³⁺ PL lifetimes

PL decay traces

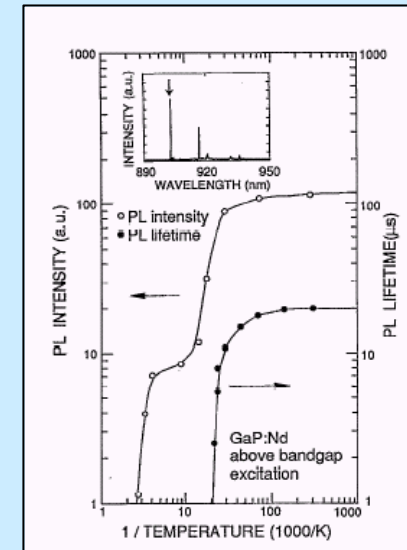


- **(A) Intrinsic nc-Si PL: $\exp(-W_{\text{nc-Si}} t)$**
 - $W_{\text{nc-Si}} < 10 \mu\text{s}$: system response due to low temp anneal (950 °C)
(*c.f.* $\sim 100 \mu\text{s}$ for high quality nc-Si after 1100 °C anneal)
- **(B) superimposition of Nd³⁺ on nc-Si PL**
 - $I(\lambda=0.92 \mu\text{m}) \sim I_{\text{nc-Si}} \exp(-t/t_{\text{nc-Si}}) + I_{\text{Nd}} \exp(-t/t_{\text{Nd}})$

Wide bandgap effect on suppression still valid?

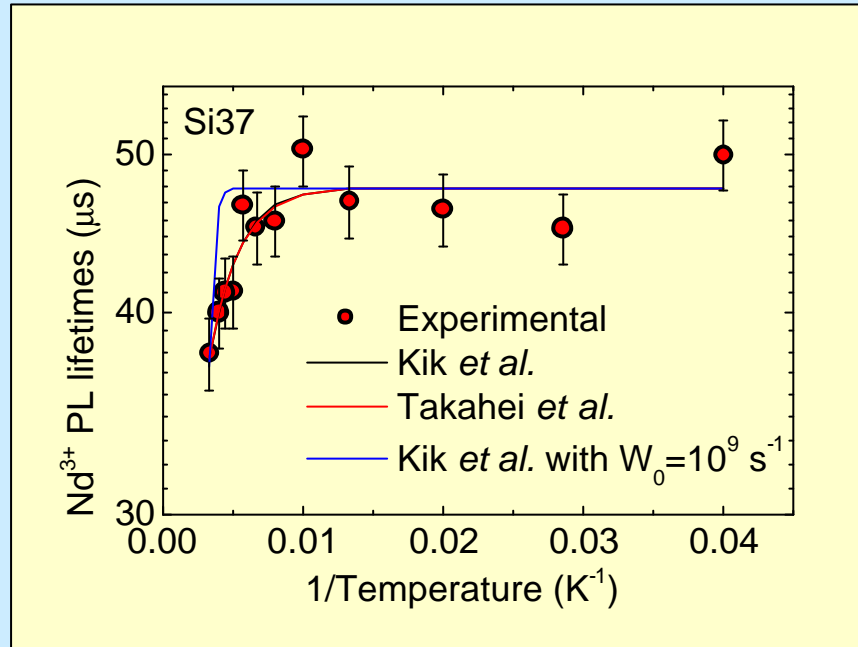


GaP:Nd (Taniguchi et al. APL, 1991)



- The lack of backtransfer
 - $t_{\text{Nd}}(T) 50 \rightarrow 40 \mu\text{s}$ with T (cf. Nd-doped silica-based thin films)
 - The reason of 3 fold decrease of Nd PL intensity
 - not backtransfer but exciton dissociation
- Only ~ 200 meV difference between nc-Si luminescence and Nd³⁺ excitation energy
 - Not sufficient to suppress backtransfer (cf. bulk GaAs:Nd)
 - Even GaP ($E_g=2.26$ eV) shows > 10 thermal-quenching of Nd³⁺ luminescence

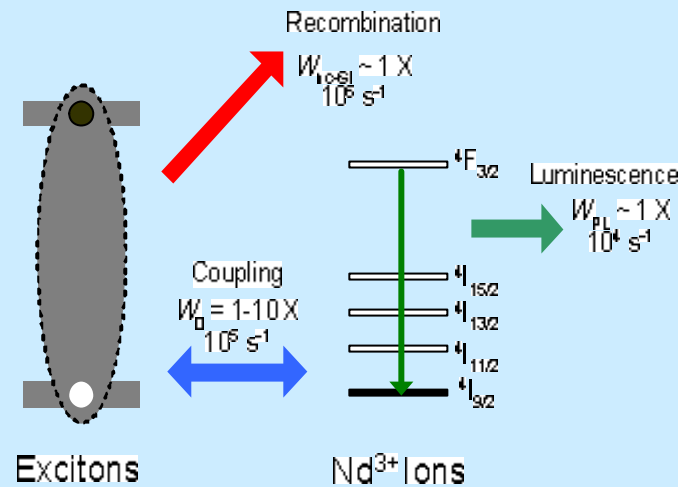
Weak couplings induce weak backtransfer



- Kik's
 - $W_{BT}(T) = W_0 \exp(-E_0/kT)$
- Takahei's
 - $W_{BT}(T) = W_0 (n_q)^p \exp(-2n_q s)$
 $n_q = [\exp(\hbar\omega/kT) - 1]^{-1}$

- Analysis of temperature quenching of with
 - Kik's (Thermal dissociated) and Takahei's (Multi-phonon relaxed) model
 - Different physical basis but similar dependence on E_0 and W_0
→ $W_{BT}(T) = W_0 \times f(E_0, T)$ (W_0 : coupling prefactor)
- Identical simulated results fit by both models with similar fitting value
 - $E_0 \sim 100 \text{ meV}$ and $W_0 \sim 1-10 \times 10^5 \text{ sec}^{-1}$
 - Not explainable with $W_0 \sim 10^9 \text{ sec}^{-1}$ (for Er-doped bulk Si)

Efficient excitation with suppression of backtransfer



- Small energy mismatch but weak coupling

- $E_0 \sim 100 \text{ meV}$: agreement with previous results (Taguchi *et al.* JAP, 1996)
- $W_0 \sim 1-10 \times 10^5 \text{ sec}^{-1}$: much smaller than $10^8-10^{10} \text{ sec}^{-1}$ (for GaAs or GaP)
 - responsible for the lack of backtransfer
 - consistent with the excitation rate of Er³⁺ 10^6 sec^{-1} (Watanabe *et al.* 2001)
- Slight separation of RE from nc-Si weakens couplings suppressing backtransfer

Implications

- Nc-Si : efficient material for RE luminescence
 - High quality nc-Si by high temperature anneal : $W_{\text{nc-Si}} < W_{\text{Nd}} < W_0$
 - The excitons created via backtransfer can re-excite RE before dissociation or radiative recombination
 - Window for the efficient excitation ($W_{\text{nc-Si}} > W_0$) and the suppression of backtransfer : unique for nc-Si
 - Moreover, suppression of de-excitation of RE by Auger-excitation of free carriers (Seo *et al.* APL 1999)
 - **Efficient excitation of RE w/o all possible de-excitation process**

- Slow excitation rate? ($W_0 = 10^5 \text{ sec}^{-1}$):
 - For example, $10 \mu\text{m} \times 10 \mu\text{m} \times 1 \text{ cm}$ Er-doped waveguide based on nc-Si:Er with ~ 0.1 at % of Er
 - $\sim 100 \text{ mW}$ output power is feasible
 - Enough for the practical application

Summary and Conclusion

- **Excitons mediated Nd excitation in nc-Si**
 - The size of nc-Si: crucial for Nd excitation in SRSO:Er
- **3 fold quenching of Nd PL intensity**
 - Due to exciton dissociation
- **The lack of backtransfer**
 - Small energy mismatch (~ 200 meV)
 - However weak couplings between nc-Si and RE
 - responsible for site dependent quenching behavior
- **Efficient RE excitation avoiding thermal quenching possible**